Boussinesq thermal convection in a rotating spherical shell with an outer stably stratified layer

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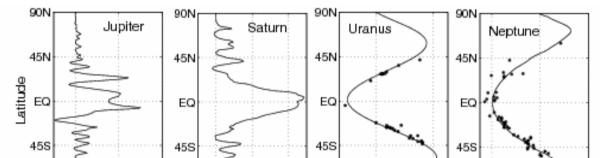
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Abstract

Finite amplitude thermal convection in a rapidly rotating spherical shell associated with a stably stratified layer placed near the outer surface is investigated. Systematic numerical experiments are performed with an Ekman number of $E=10^{-3}$, a Prandtl number of P=1 and an inner/outer radius ratio of $\eta=0.4$, and the existence of a strongly stratified upper layer is shown to enhance the generation of equatorial surface retrograde flows when the Rayleigh number is approximately ten times larger than the critical value. The existence of the stable layer causes the bottom of the stable layer to behave as a virtual boundary for the convective motion underneath. Its effective dynamic condition varies from the free-slip condition to the no-slip condition as the Rayleigh number increases. Budget analysis of angular momentum shows that the Reynolds stress of the convective vortices beneath the stable layer is weakened and is dominated by the transport of the planetary angular momentum when the Rayleigh number is increased. As a result, the latitudinal temperature gradient produced at the bottom of the stable layer induces the equatorial retrograde flow through the thermal wind balance. This diffuses through the stable layer by viscosity and produces the equatorial surface retrograde flow.

1. Introduction

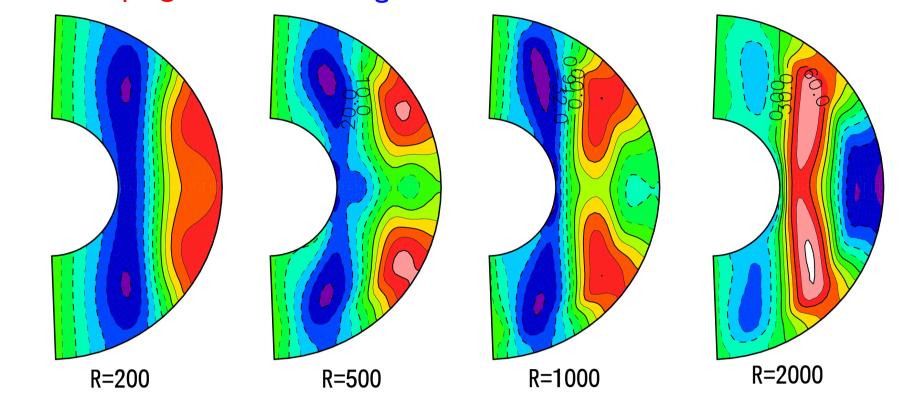
OGas giant planets : equatorial superrotation



3. Results

OMean zonal flows : Rayleigh number dependency

Red: prograde. Blue: retrograde. The cases with Γ =100.

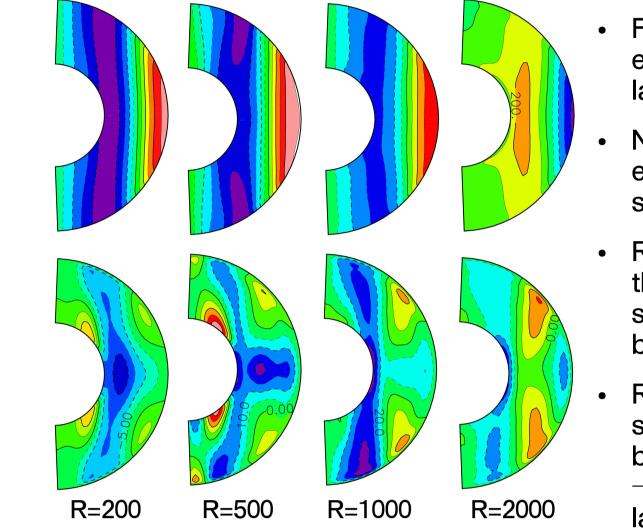


4. Discussions

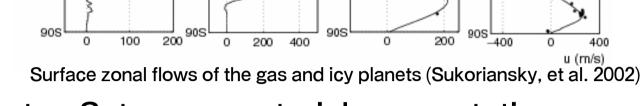
OExperiments with a boundary just beneath the stable layer

Red:prograde. Blue: retrograde.

Upper : free-slip at the top. Lower : no-slip at the top.



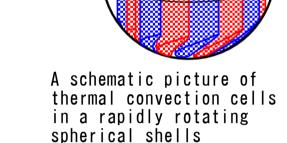




- Jupter, Saturn : equatorial superrotation
- Equatorial superrotation : dynamically interesting
 - Simple angular mom.transport \rightarrow equatorial surface retrograde flow
 - Some acceleration mechanisms needed

OThermal convection in rotating spherical shells

- A model for explaining equatorial superrotaion
- In rapidly rotating spherical shells
 - \rightarrow Taylor-column type convection uniform along the rotating axis
- prograde- outward tilting of the convection cells
 - \rightarrow momentum transport by Reynolds stress
 - \rightarrow Equatorial superrotation
- When the Rayleigh number is increased \rightarrow mixing of angular mom.
 - \rightarrow retrograde equatorial flows



$\bigcirc A$ stably stratified layer near the top boundary?

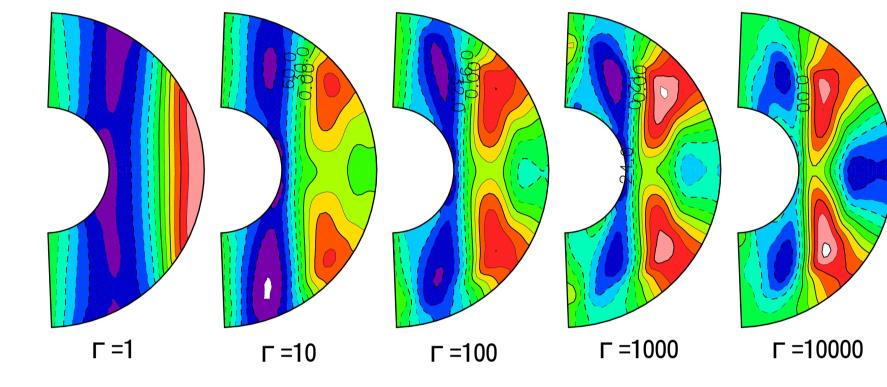
- Most of studies : the whole layer is unstable
- Stably stratified layers in the upper part of the actual planetary atmospheres (stratospheres, tropospheres)
- The existence of stable layers inhibits columnar convection \rightarrow angular mom.transport possible?

- Small Rayleigh numbers : equatorial prograde at the top
- Large Rayleigh numbers : equatorial retrograde at the top

\bigcirc Mean zonal flows :

dependence of the stable stratification

Red: prograde. Blue: retrograde. The cases with R=1000.



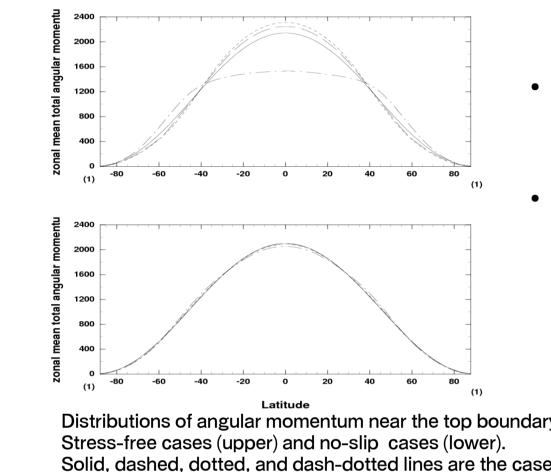
- Weak stratification : equatorial prograde at the top
- Strong stratification : equatorial retrograde at the top

ODirection of equatirial surface zonal flows

equatorial retrograde at large Rayleigh numbers

- No-slip at the top : equatorial retrograde at small Rayleigh numbers
- R≧500: the distributions are similar to the no-slip boundary cases
- R=200 : similar to the free-slip boundary cases \rightarrow the effect of the stable layer is a medium condition between no-slip and free-slip

ODistribution of angular momentum

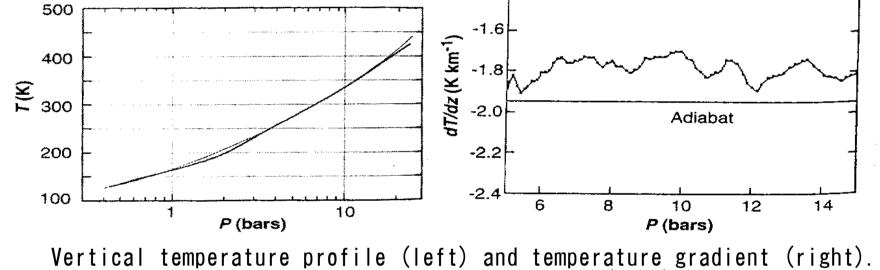


• **Retrograde** with free-slip boundary ← angular mom.homogenization

Retrograde with no-slip boundary : inhomogeneous angular mom. \leftarrow weakening of Reynolds stresses by viscosity?

Distributions of angular momentum near the top boundary. Solid, dashed, dotted, and dash-dotted lines are the cases with $R = 2 imes 10^2, 5 imes 10^2, 1 imes 10^3, 2 imes 10^3,$ respectively.

OAngular momentum budget analysis



Note that stable stratification. Seiff et al. (1996)

OPrevious studies

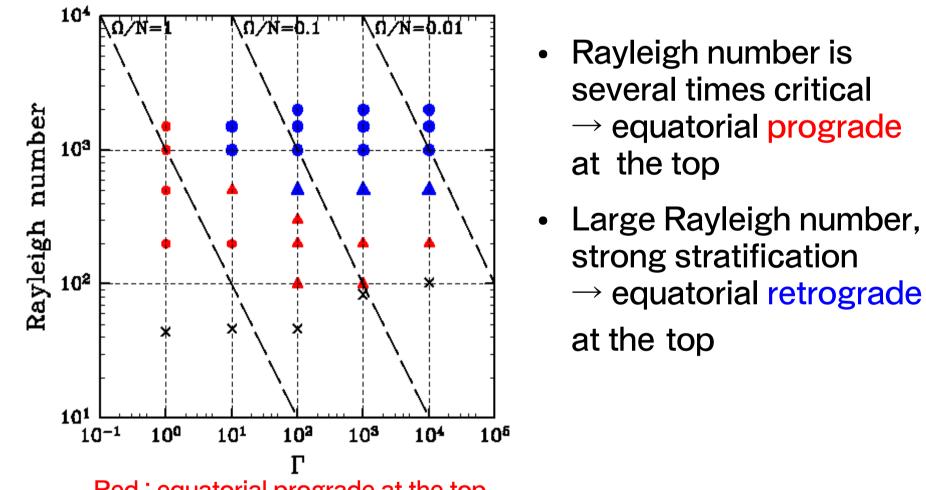
- Zhang and Schubert (1996,1997) Critical convection. Columnar convection completely penetrating a stable layer
- Takehiro and Lister (2001) The penetration thickness is in proportion to the ratio between the angular velocity and Brunt-Vaisala frequency Ω/N
- Takehiro and Lister (2002)

Finite amplitude convection with the Rayleigh numbers several times of the critical. Mean zonal flows penetrates by viscosity, resluting equatorial superrotation.

OThis study

- The dependency of equatorial surface mean zonal flows on the Rayleigh number and strength of stably stratified layers?
- The Rayleigh number is further increased.
- The stability of stable layer is varied more widely.

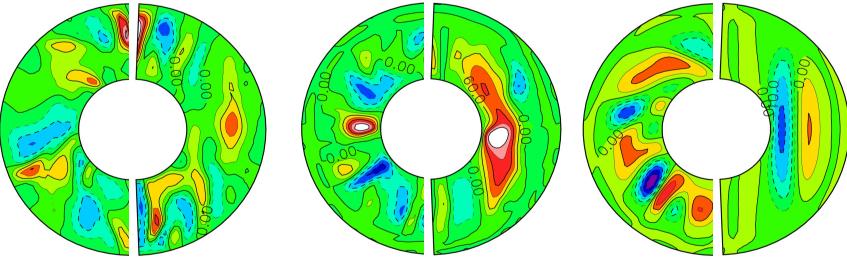
Model



Red : equatorial prograde at the top. Blue: equatorial retrograde at the top.

\bigcirc Convective flows

Left : equatorial cross section. Right : meridonal cross section Red : upward flows. Blue : downward flows

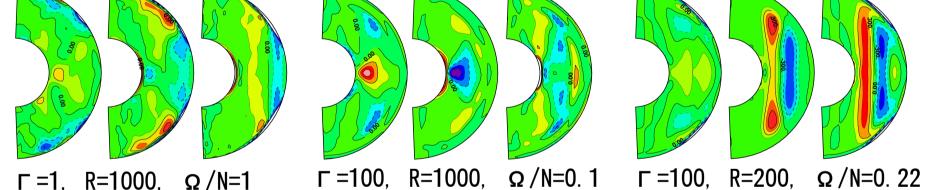


 $\Gamma = 100, R = 1000, \Omega / N = 0.1 \Gamma = 100, R = 200, \Omega / N = 0.22$ Γ=1, R=1000, $\Omega/N=1$

- Weak stratification : convective motions penetrate the stable layer
- Strong stratification, large Rayleigh numbers : trapped below the stable layer
- Strong stratification, small Rayleigh numbers : columnar convection below the stable layer

• Conservation of angular momentum $\frac{\partial}{\partial t}(\overline{u}_{\phi}r\sin\theta) + \nabla\cdot\mathbf{F}_{p} + \nabla\cdot\mathbf{F}_{r} + \nabla\cdot\mathbf{F}_{d} + \nabla\cdot\mathbf{F}_{v} = 0.$ $\mathbf{F}_p = \frac{1}{E} r^2 \sin^2 \theta (\overline{u}_\theta \mathbf{e}_\theta + \overline{u}_r \mathbf{e}_r)$ Planetary angular mom. flux $\mathbf{F}_r = \overline{u}_{\phi} \overline{u}_{\theta} r \sin \theta \mathbf{e}_{\theta} + \overline{u}_{\phi} \overline{u}_r r \sin \theta \mathbf{e}_r$ Relative angular mom. flux $\mathbf{F}_d = \overline{u'_{\phi}u'_{\theta}}r\sin\theta\mathbf{e}_{\theta} + \overline{u'_{\phi}u'_r}r\sin\theta\mathbf{e}_r$ Reynolds stresses $\mathbf{F}_{v} = -r^{2}\sin^{2}\theta\nabla\left(\frac{\overline{u}_{\phi}\sin\theta}{r^{2}\sin^{2}\theta}\right)$ Viscous angular mom. flux

Distributions of zonally and temporally averaged angular momentum flux divergences. Left, center, and right panels are $\nabla \cdot \mathbf{F}_p, \nabla \cdot \mathbf{F}_d, \nabla \cdot \mathbf{F}_p + \nabla \cdot \mathbf{F}_r + \nabla \cdot \mathbf{F}_d$, respectively. Red : divergence. Blue : convergence.



- Weak stratification, large Rayleigh numbers: Reynolds stress dominates \rightarrow equatorial prograde
- Strong stratification, large Rayleigh numbers : Planetary angular momentum flux dominates \rightarrow equatorial retrograde
- Strong stratification, small Rayleigh numbers : Reynolds stress dominates \rightarrow equatorial prograde

5. Summary

- Large Rayleigh numbers, strong stratification \rightarrow equatorial retrograde flows
- Deep convection is turbulent rather than regular and columnar, but does not penetrate the stable layer

• System : Boussinesq fluid in a rotating spherical shell

$$E\left(\frac{\partial \boldsymbol{u}}{\partial t} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} - \nabla^2 \boldsymbol{u}\right) + 2\boldsymbol{k} \times \boldsymbol{u} + \nabla p = R\frac{\boldsymbol{r}}{r_o}T,$$
$$\frac{\partial T}{\partial t} + \boldsymbol{u} \cdot \nabla T + u_r \cdot \frac{dT_B}{dr} = \frac{1}{P}\nabla^2 T, \quad \nabla \cdot \boldsymbol{u} = 0.$$

- Parameters
 - The Rayleigh number: R
 - The Ekman number: $E = 10^{-3}$
 - The Prandtl number: P = 1
 - The ratio of the inner and outer radii: $\eta = 0.4$
 - The basic temperature (Stability of the stable layer): $T_B(r)$

OExperimental setup

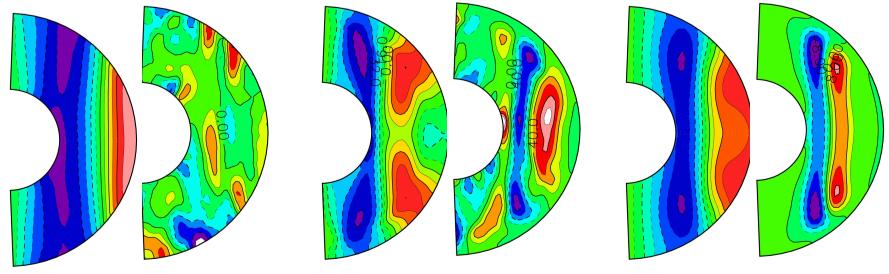
• The basic temperature profile

 $\frac{dT_B}{dr} = -\frac{1}{2}(r+\Gamma) \left| 1 - \tanh \frac{r-r_b}{a} \right| + \Gamma$

- The outer part : stable, constant temp.grad.: \prod
- The inner part : unstable, homogeneous heat source
- The boundary conditions :
 - fixed temperature, stress free
- The parameters for experimens
 - The Rayleigh number : several to 10 times critical
 - Stable stratification : $\Gamma=1\sim 10^4$

\bigcirc Mean zonal flows and disturbances

Left: mean zonal flows. Right : disturbances. Red : prograde. Blue : retrograde.



- $\Gamma = 100, R = 1000, \Omega/N = 0.1 \Gamma = 100, R = 200, \Omega/N = 0.22$ **□ =1**, R=1000. Ω/N=1
- Weak stratification : disturbances erode the stable layer
- Strong stratification : disturbances are trapped below the stable layer

- The existence of the stable layer
 - Changes the radius ratio of the convective region
 - Boundary condition at the top of the convective region \rightarrow free-slip to no-slip as the Rayleigh number is increased
 - Reynolds stress is weakened and transport by mean meridional circulation dominates

<u>References</u>

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r=1.5 r=1.67

Stable layer transition layer

unstable

r=0.67