

一様抵抗MHDにおける プラズモイド不安定性 (PI) は物理か？ MHD study of Plasmoid Instability (PI) with uniform resistivity

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Plasmoid Instability (PI)

- Multiple reconnection process.
- For high S (Lundquist number), i.e. weak uniform resistivity η , the reconnection process can become extremely bursty and stochastic, generating high reconnection rate ηJ independent of S ($\propto Rm \propto 1/\eta$).
- In these ten years, PI became very popular. Many MHD studies (Bhattacharjee, et.al., Loureiro, et.al., Cassak, et.al., Shibayama, et.al.,,,,) have been published.

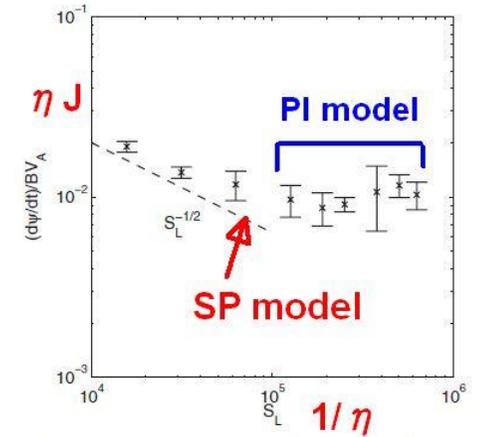
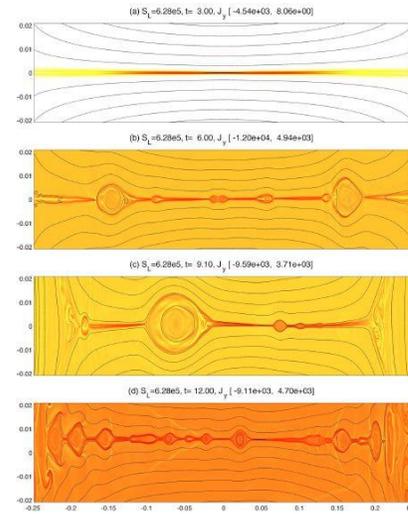


FIG. 3. Averaged reconnection rate as a function of S_L . Below $S_L \approx 3 \times 10^4$, the reconnection rate is Sweet-Parker and decays as $S_L^{-1/2}$, shown by the dotted line. Above this threshold, the reconnection rate becomes weakly dependent on S_L .

Bhattacharjee, et.al, PoP2009

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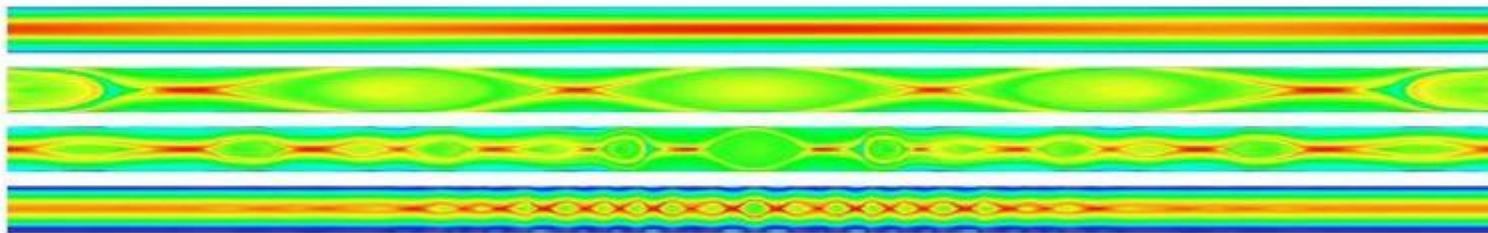
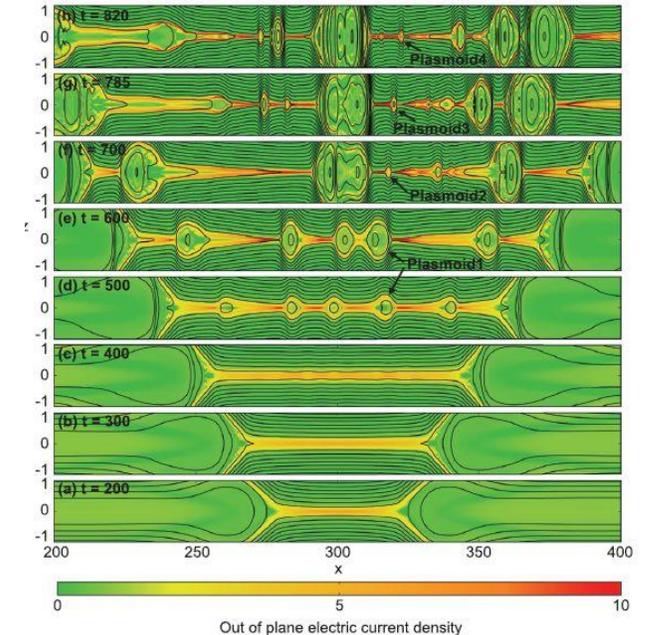


FIG. 2 (color online). Current density for $S = 10^4$, $S = 10^5$, $S = 10^6$ and $S = 10^7$. $S = 10^8$ is shown in Fig. 1.

Loureiro, et.al, PhysRevLett2009



Shibayama, et.al, PoP2015

History of MHD reconnection process

<< Theoretical MHD models >>

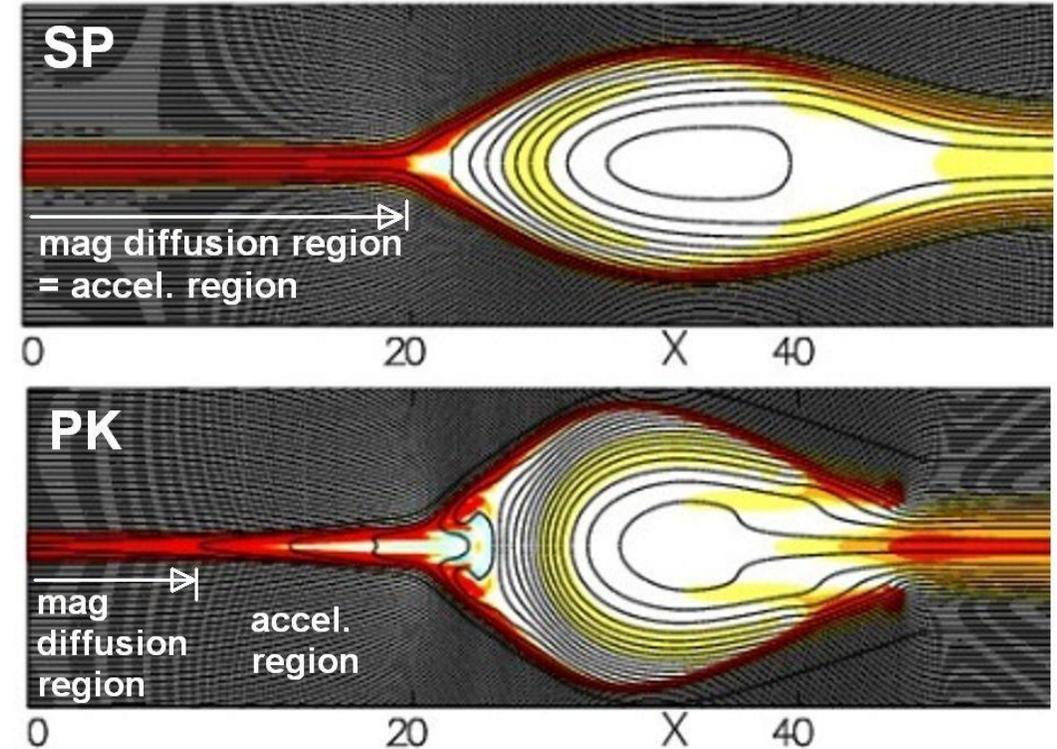
- **Sweet-Parker (SP)** model (Sweet 1958 & Parker) →
- **Petschek (PK)** model (Petschek, 1964) →
- **FKR theory** (1963) (linear tearing instability)

<< Numerical MHD studies >>

- Spontaneous model (Ugai, 1977) vs Externally driven model (Sato, 1978)
- **Uniform resistivity** model vs **Anomalous** (non-uniform) resistivity model

<< Our Final Goal is “What should be the “fast” mag. rec. ?” >>

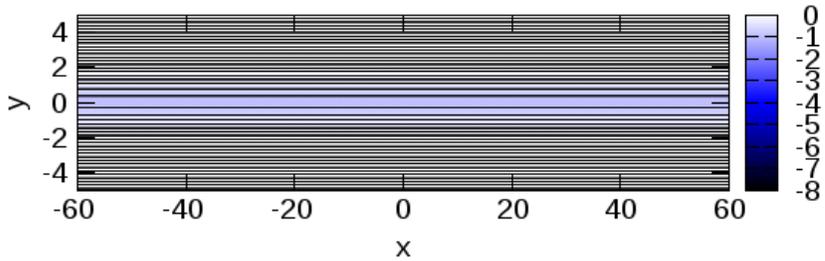
- Spontaneous & Anomalous resistivity → **PK model**
- Spontaneous, Uniform resistivity & FKR → **PI model** (= multiple SP model)



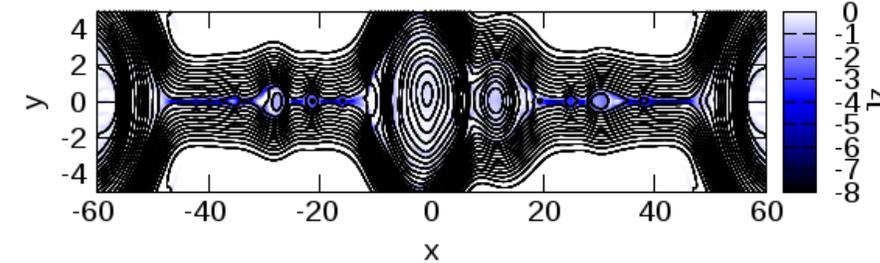
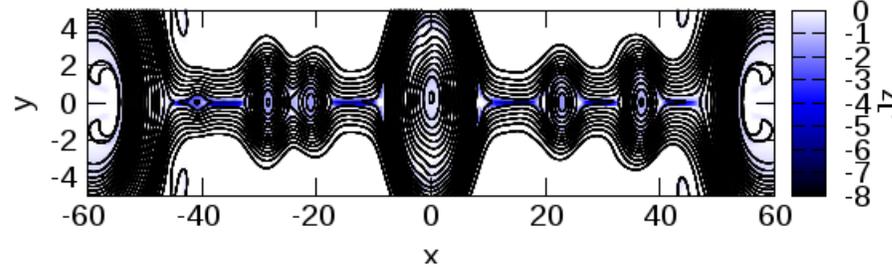
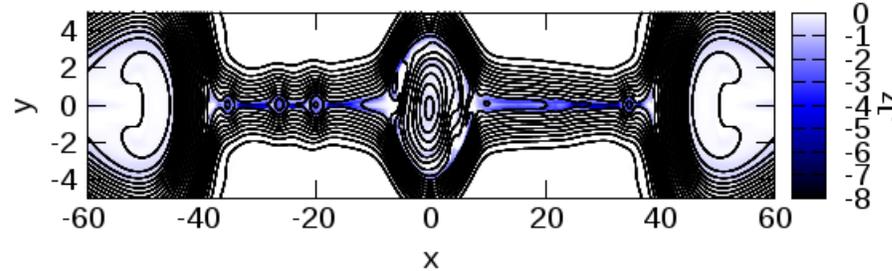
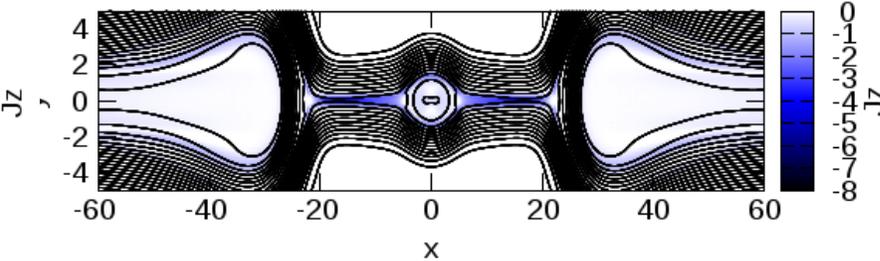
HLLD numerical scheme

- **Excellent shock capturing scheme** for MHD models
- Region size $(L_x, L_y) = 100 * 20$
- Grid size $(N_x, N_y) = 3000 * 600 \sim 12000 * 2400$ (e.g., $N_x = 3000$ is $dx = dy = 1/15$)
- HLLD setup
 - Slope limiter :minmod
 - Time marching :TVD RK2nd
 - Time step is automatically adjusted by CFL condition as time proceeds.
- Initial temperature is uniform without guide field. Initial upstream is Beta=0.15.
- Every boundary condition is symmetry, i.e. closed system model.
- For $0 < T < 5$, a resistive disturbance to initiate mag rec is locally put around origin.
- For $5 < T$, the resistive disturbance is removed. Instead, uniform resistivity is maintained. Then, wait for PI.

$T=130 \sim 140$ for $Rm=50, 100, 250, 1000, \text{ and } \infty$

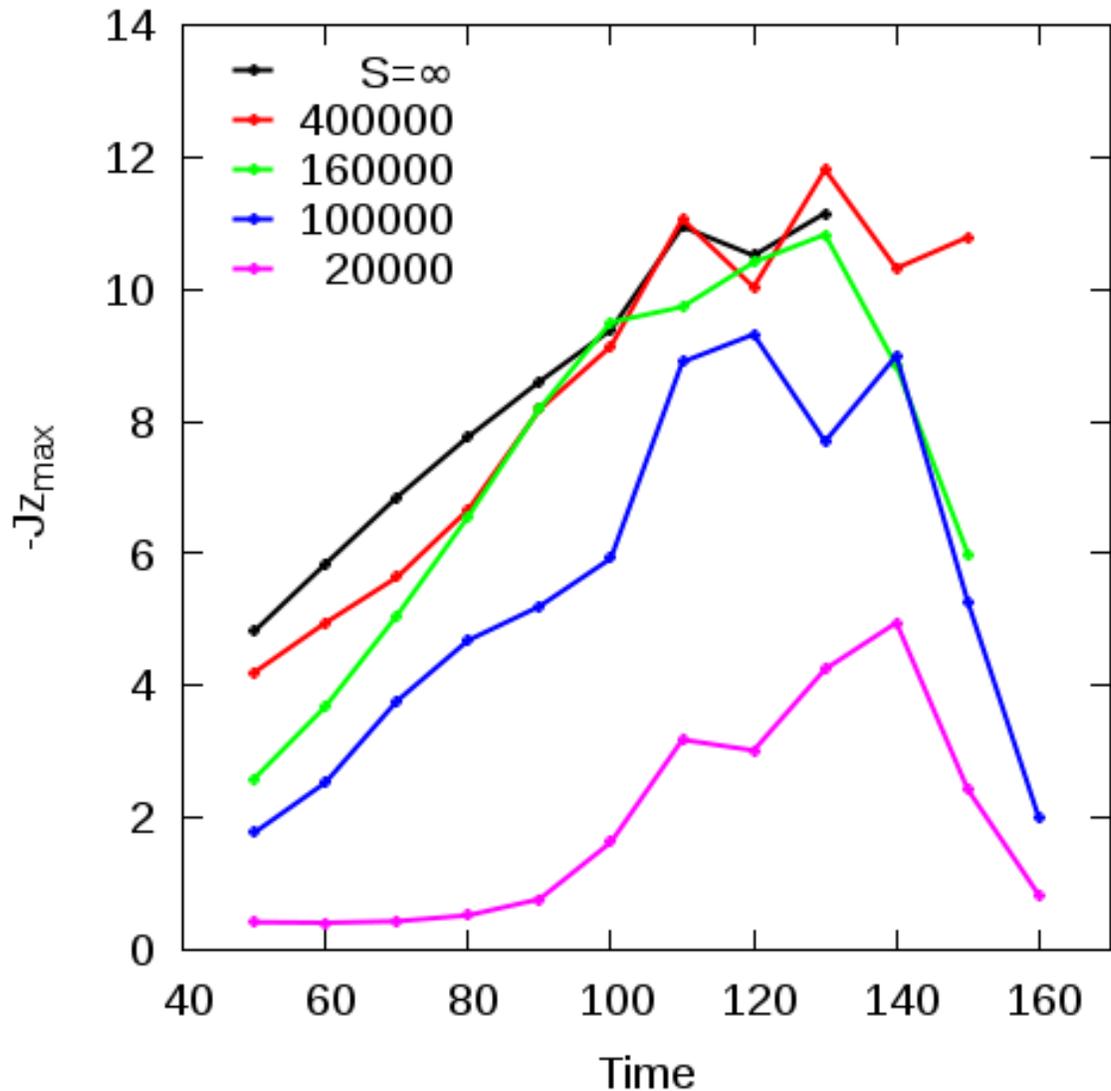


1D current sheet at $T=0$



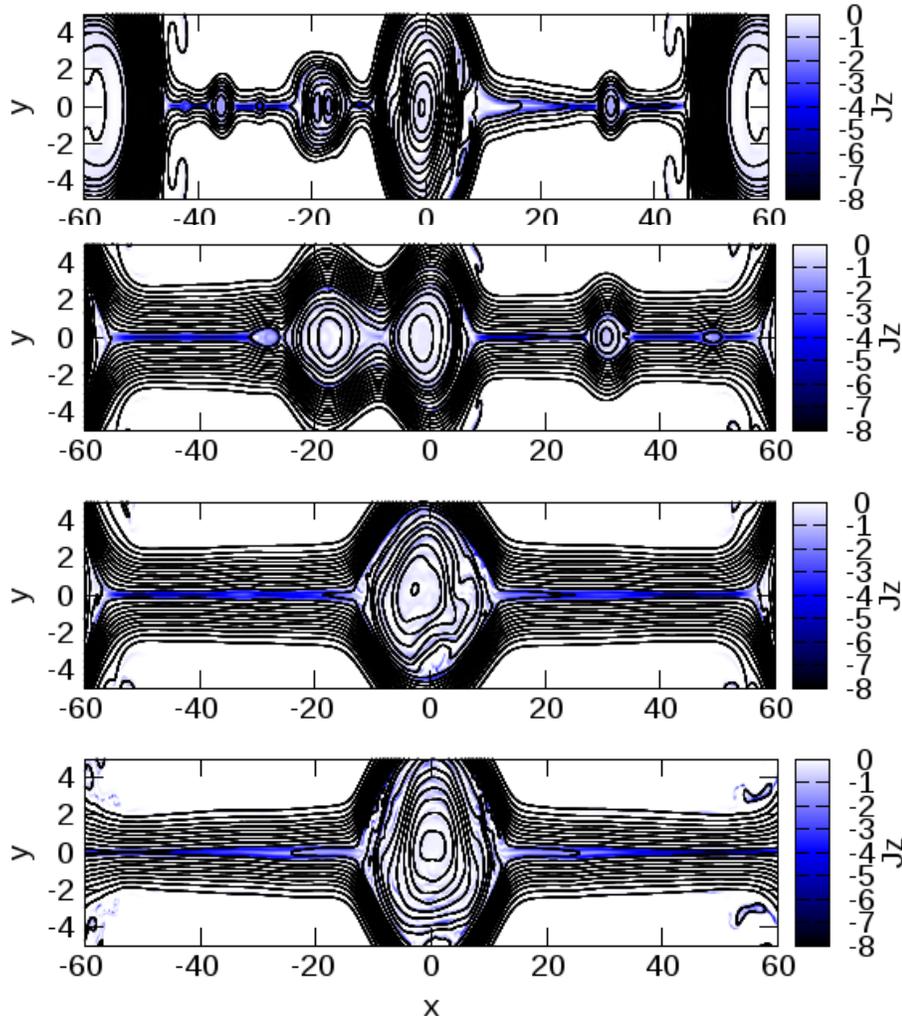
- ランキスト数 S (磁気レイノルズ数 Rm) を上げるとテアリングが多発する。
- ただし、 $S(Rm)$ 無限大でもテアリングが起こるのは明らかに**数值的効果** (計算誤差)。

Time variations of J_z at X-point are numerically saturated.



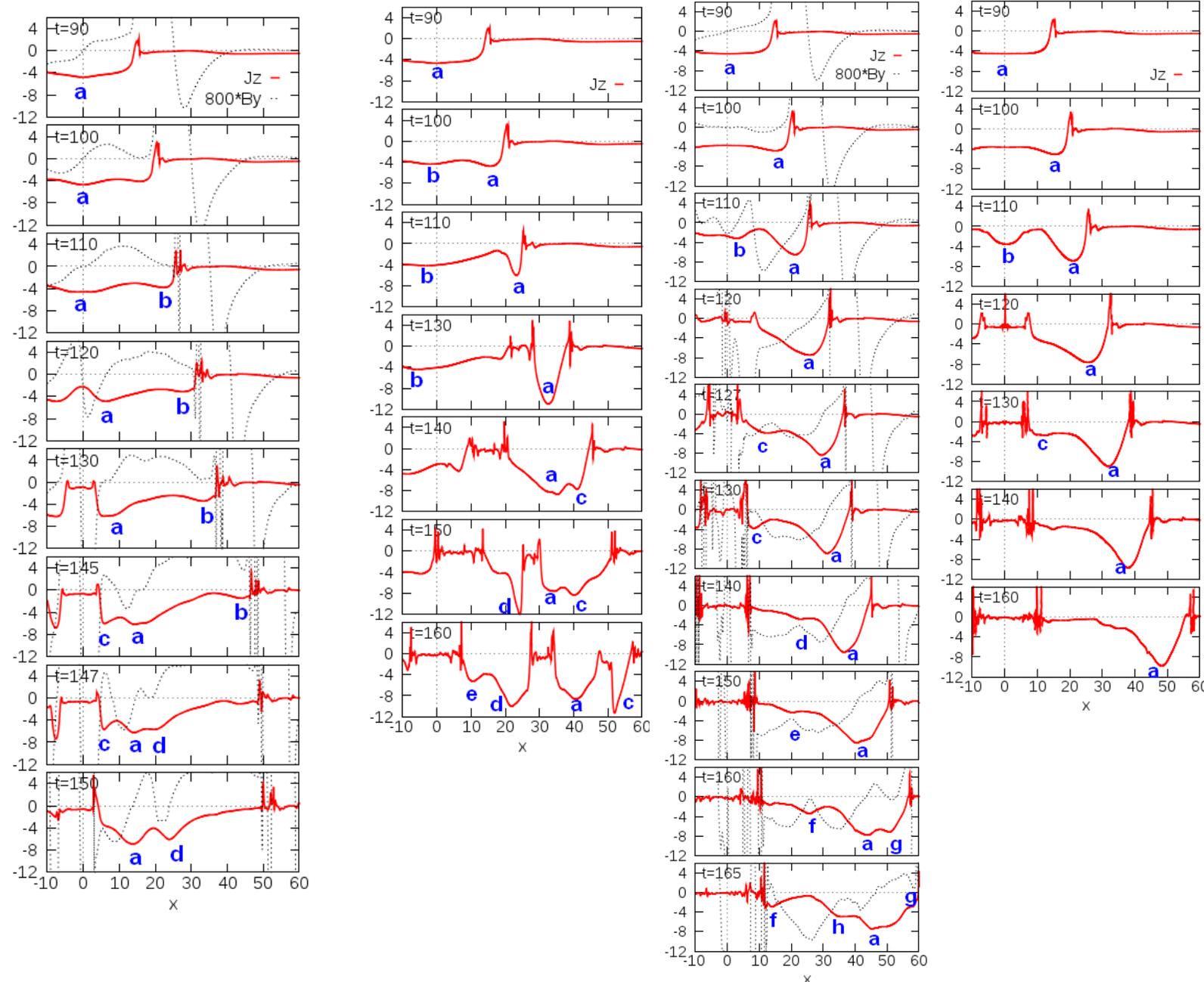
- J_z -plot at the most active (highest J_z) X-point.
- $R_m=\infty$ ($S=\infty$)
- $R_m=1000$ ($S=400000$)
- $R_m=400$ ($S=160000$)
- $R_m=250$ ($S=100000$)
- $R_m=50$ ($S=20000$)
- $R_m>400$ は J_z の時間変化が数値的に飽和している。

As $1/dx$ increases from 15 to 120 for $Rm=250$, the plasmoid instability (PI) becomes less active.



- $dx=1/15$, $T=140$
- $dx=1/45$, $T=160$
- $dx=1/60$, $T=160$
- $dx=1/120$, $T=170$
- 数値解像度 $1/dx$ を上げると
PIは徐々に消える。

Time variations of Jz-profiles on x-axis



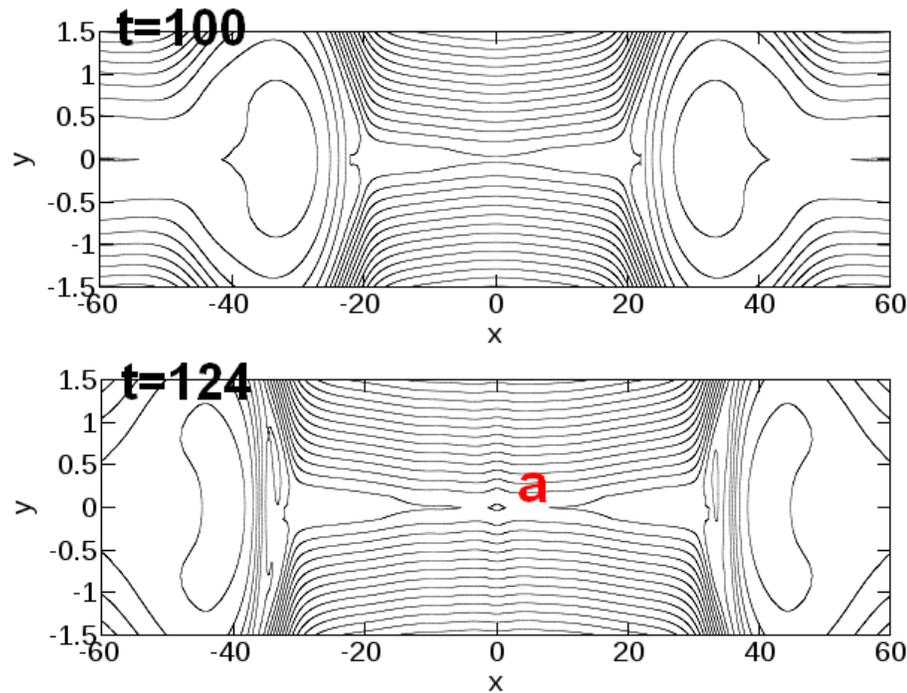
- $dx=1/30, 1/45, 1/60, 1/90$ の図
- ピーク ($J_z < 0$) の高さ順にアルファベットで命名。
- $dx=1/30$ から $1/45$ への変化では、途中でJzピークaとb(大小)が入れ替わる(相転移)。
- つまり、数値誤差によりある種の相転移が起こっている。
- $dx=1/45$ から $1/90$ への変化はPIがless activeになる傾向。

A phase transition (drastic switching of Jz-peaks "a" and "b") is observed between 1/30 and 1/45, where alphabetical order means the order of higher Jz-peaks.

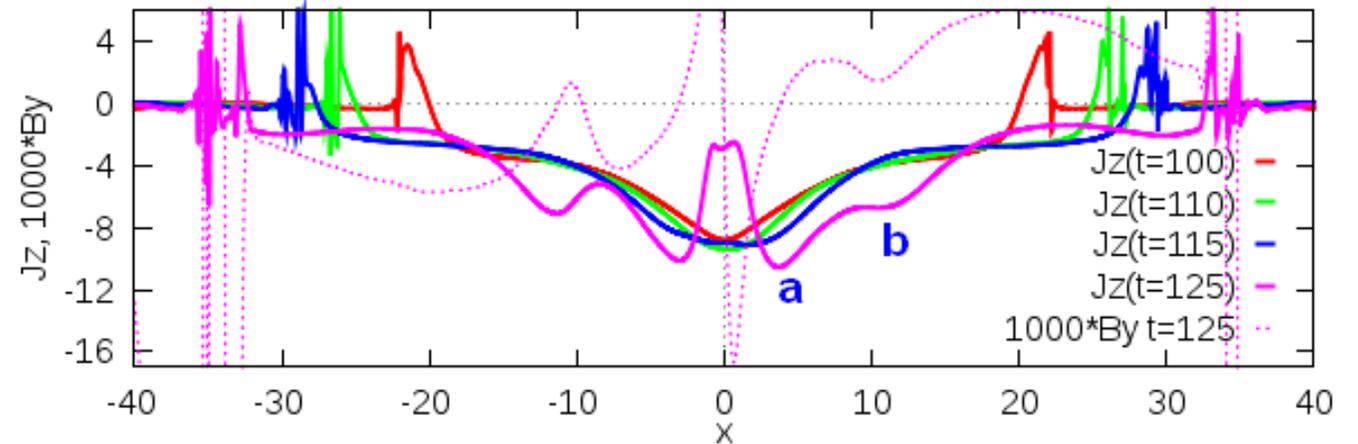
Numerical converging observed from 1/45 to 1/90 shows PI tends to be monotonically less active.

対称テアリングと非対称テアリング

symmetric and asymmetric tearing instability (exactly, two types of J_z -peak separations)

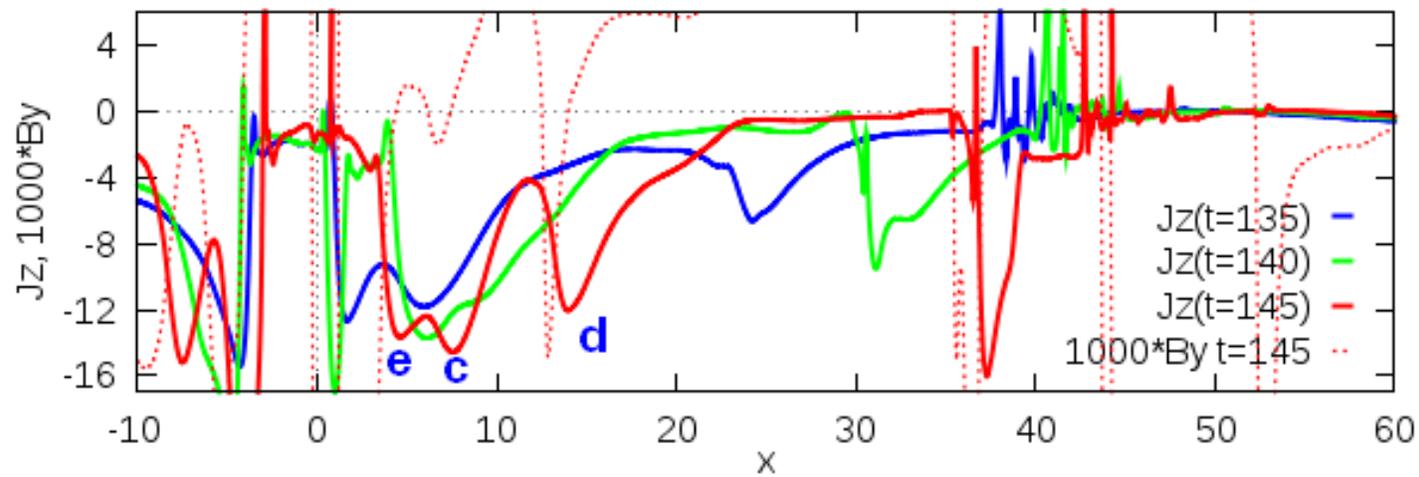
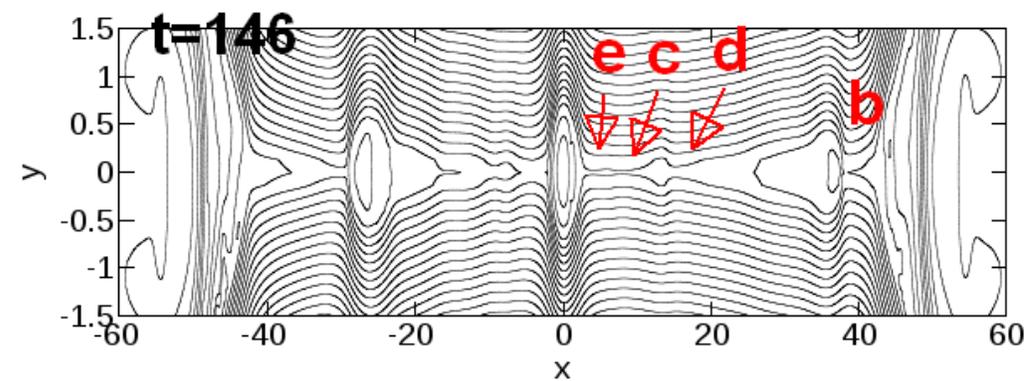
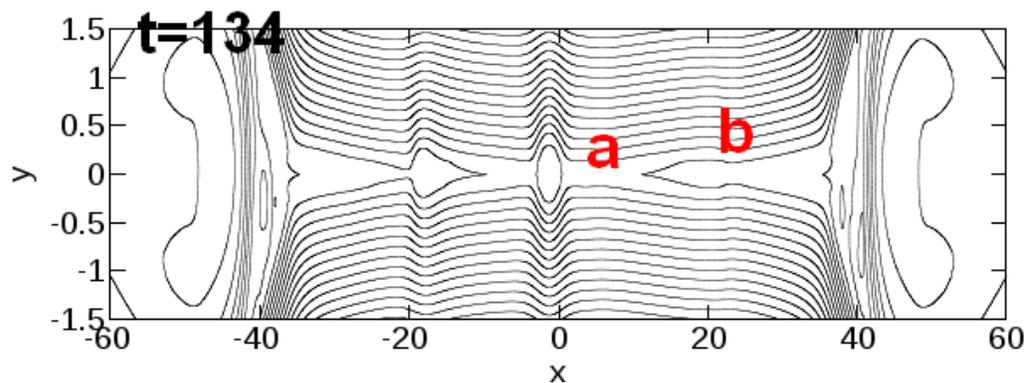


$Rm=1000, dx=/90, 100 < T < 125$

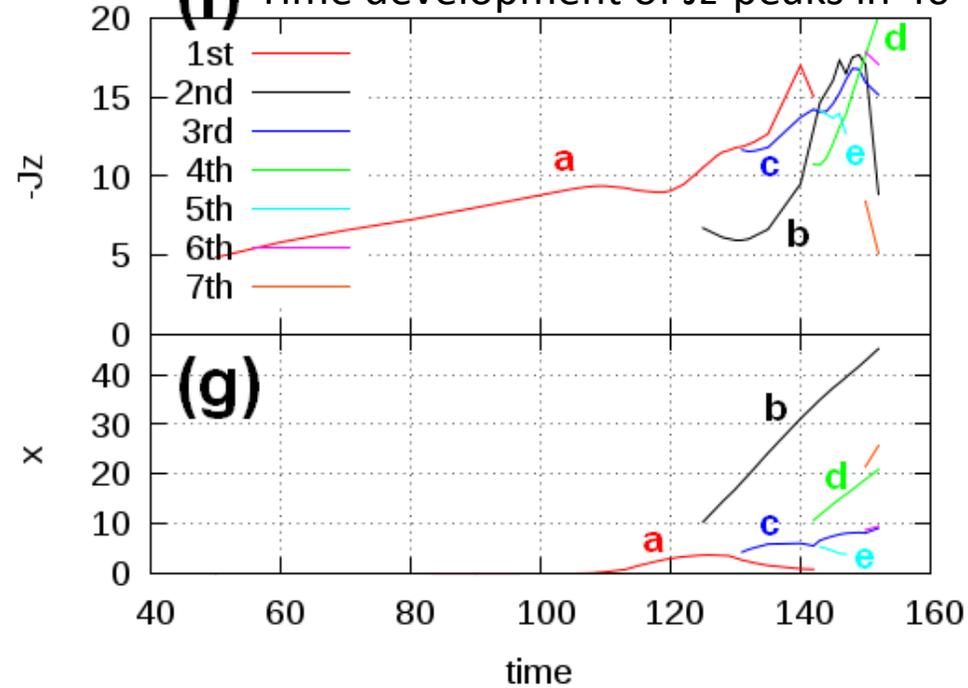


- Observe that a J_z -peak is bifurcated to two J_z -peaks.
- The bifurcation can be classified into symmetric (a) and asymmetric (b) case.
- The symmetric case (a) tends to affect a phase transition for larger $1/dx$.

Rm=1000, dx=1/90, 134<T<146



(f) Time development of J_z -peaks in $40 < T < 150$



まとめ

- 一様抵抗MHDで起こるPIは数値的な影響を受けて速い(PKライクな)リコネクションになる。
- 数値的な影響を排除すると、遅い(SPライクな)リコネクションになる。
- MHD-PI with uniform resistivity is studied with HLLD code.
- As numerical resolution $1/dx$ increases, PI tends to be less active where drastic Jz-peak separations may be caused as phase transitions.
- Symmetric tearing instability, i.e. symmetric Jz-peak separation tends to affect the phase transition due to numerical errors.
- We cannot numerically confirm that the fully developed PI is reproduced. It suggests that Sc (critical Lundquist number) does not exist.