Disappearance of surface banded structure produced by thermal convection in a rapidly rotating thin spherical shell

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Japanese-French model studies of planetary atmospheres @ CPS, Kobe Univ.





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2 Numerical experiments

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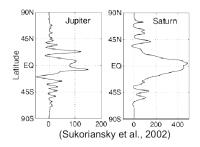


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## Surface flows of gas giant planets

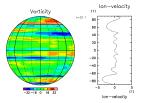
- Surface flows of Jupiter and Saturn are characterized by the broad prograde jets around the equator and the narrow alternating jets in mid- and high-latitudes.
- It is not yet clear whether those surface jets are produced by convective motions in the "deep" region, or are the result of fluid motions in the "shallow" weather layer.

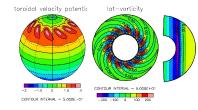




# "Deep" models and "Shallow" models

- "Shallow" models
  - 2D turbulence on a rotating sphere
  - Primitive model
    - Result: Narrow alternating jets in mid- and high-latitudes.
    - Problem: the equatorial jets are not necessarily prograde
  - Recent improvements
    - Strong Newtonian cooling (shallow water)
    - MHD drag (primitive model)
- "Deep" models
  - Convection in rotating spherical shells
    - Result: Produce equatorial prograde flows easily
    - Problem: difficult to generate alternating jets in mid- and high-latitudes

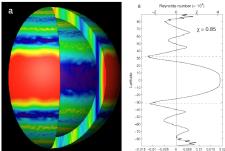




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# "Thin" spherical shell model

- Heimpel and Aurnou (2007) (hereafter, HA2007)
  - "Thin" spherical shell model with large Rayleigh number, small Ekman number.
  - Prograde jets and alternating jets in midand high-latitudes can produce simultaneously
  - However, eight-fold symmetry in the longitudinal direction is assumed.



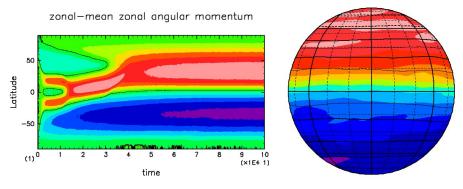
- HA2007 explains that the banded structure in high latitudes is caused by forced 2D turbulence on the  $\beta$  plane (Rhiens effect)
- Forced 2D barotropic turbulence on a rotating sphere : banded structure disappears after long time integration (Obuse et al. 2010).
- Question : does the banded structure of HA2007 disappear after long time integration or not?

#### In the present study:

- Let us perform long time integration of thermal convection in a rotating thin spherical shell.
- Calculate in the whole domain.

## Forced 2D barotropic turbulence on a rotating sphere

- Obuse et al. (2010)
  - $\bullet\,$  Merger and disappearance of jets  $\Rightarrow$  banded structure disappears



Time development of latitudinal distribution of angular momentum (left) and final zonal wind profile (right) (Obuse et al. 2010)



2 Numerical experiments

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• Boussinesq fluid in a rotating spherical shell.

• Scaling: the shell thickness, rotation period, temperature difference.

$$\nabla \cdot \boldsymbol{u} = 0,$$
  

$$\frac{\partial \boldsymbol{u}}{\partial t} + (\boldsymbol{u} \cdot \nabla)\boldsymbol{u} - E\nabla^{2}\boldsymbol{u} + 2\boldsymbol{k} \times \boldsymbol{u} + \nabla p = \frac{\operatorname{Ra}E^{2}}{\operatorname{Pr}} \frac{\boldsymbol{r}}{r_{o}} T,$$
  

$$\frac{\partial T}{\partial t} + (\boldsymbol{u} \cdot \nabla)T = \frac{\operatorname{E}}{\operatorname{Pr}} \nabla^{2} T.$$
  
• Parameters:  
• Prandtl number:  $\operatorname{Pr} = \frac{\nu}{\kappa}$   
• Rayleigh number:  $\operatorname{Ra} = \frac{\alpha g_{o} \Delta T D^{3}}{\kappa \nu}$   
• Ekman number:  $\operatorname{E} = \frac{\nu}{\Omega D^{2}}$   
• radius ratio:  $\eta = \frac{r_{i}}{r_{o}}$ 

- Boundary condition: Isothermal, Impermeable and Stress free.
- Input parameters:

Prandtl number: Pr	0.1
Radius ratio: $\eta$	0.85
Ekman number: ${ m E}$	$3 \times 10^{-6}$
Modified Rayleigh number: ${\rm Ra}^{\ast}$	0.05

- the definition of modified Rayleigh number:  $\operatorname{Ra}^* = \frac{\operatorname{Ra}E^2}{\operatorname{Pr}} = \frac{\alpha g \Delta T}{\Omega^2 D}$ 
  - Ratio of Coriolis term and buoyancy term

## Numerical methods

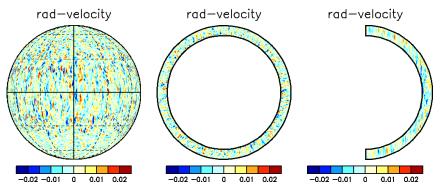
- Traditional spectral method.
  - Toroidal and Poloidal potentials of velocity are introduced.
  - The total wave number of spherical harmonics is truncated at 341, and the Chebychev polynomials are calculated up to the 48th degree.
    - The numbers of grid points:1024, 512, and 65 in the longitudinal, latitudinal, and radial directions, respectively.
- In order to save computational resources, we use hyperdiffusion with the same functional form as the previous studies

$$\nu = \left\{ \begin{array}{ll} \nu_0, & \text{for } l \leq l_0, \\ \nu [1 + \varepsilon (l - l_0)^2], & \text{for } l > l_0. \end{array} \right.$$

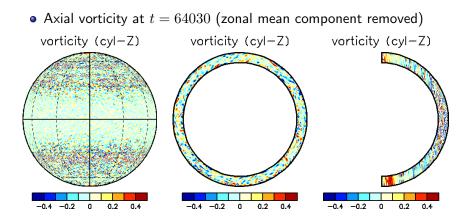
- we choose  $l_0 = 21, 42, 85, 170, \varepsilon = 10^{-2}$  ( $l_0$  is gradually increased)
- The time integration is performed using the Crank-Nicolson scheme for the diffusion terms and the second-order Adams-Bashforth scheme for the other terms.

## Convective activity

• Radial velocity at t = 64030 (about 10000 rotation).



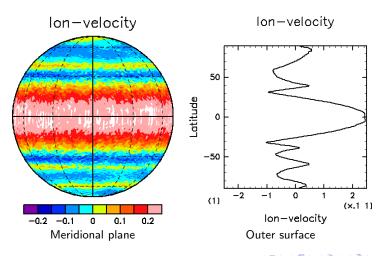
Vorticity



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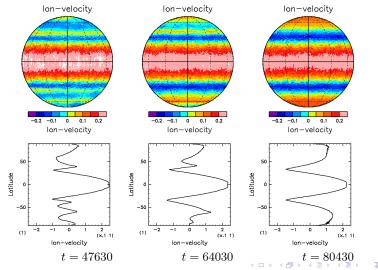
#### Mean zonal flows

- Mean zonal flow at t = 47630 (about 7500 rotation).
  - Banded structure appears in mid- and high-latitudes.



## Development of zonal flows

- Further time integration.
  - Number of jets decreases



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Thermal convection in rotating thin shells

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# Summary

- Disappearance of alternating jets in high latitudes
  - Prograde acceleration in high latitudes
  - Retrograde acceleration around the tangential cylinder
- What is acceleration mechanism? One possible story is...
  - Thermal convection occurs in high latitudes
  - It excites 2-dim. axial vortices
  - The vortices propagate as Rossby waves outward and are absorbed near the tangent cylinder
  - Associated with the propagation of Rossby waves, negative angular momentum is transferred outward
- Deep models may by difficult to explain the banded structure of Jupiter and Saturn?

#### Acknowledgement

The numerical computation of thermal convection in a rotating thin spherical shell was carried out on the Earth Simulator (ES2) at the Japan Agency for Marine Earth Science and Technology.



2 Numerical experiments

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