

固体核の形成

微惑星の成長

微惑星の成長率

$$\frac{dm}{dt} = \pi s^2 \Theta \frac{\Sigma_s}{2h_s} \delta v$$

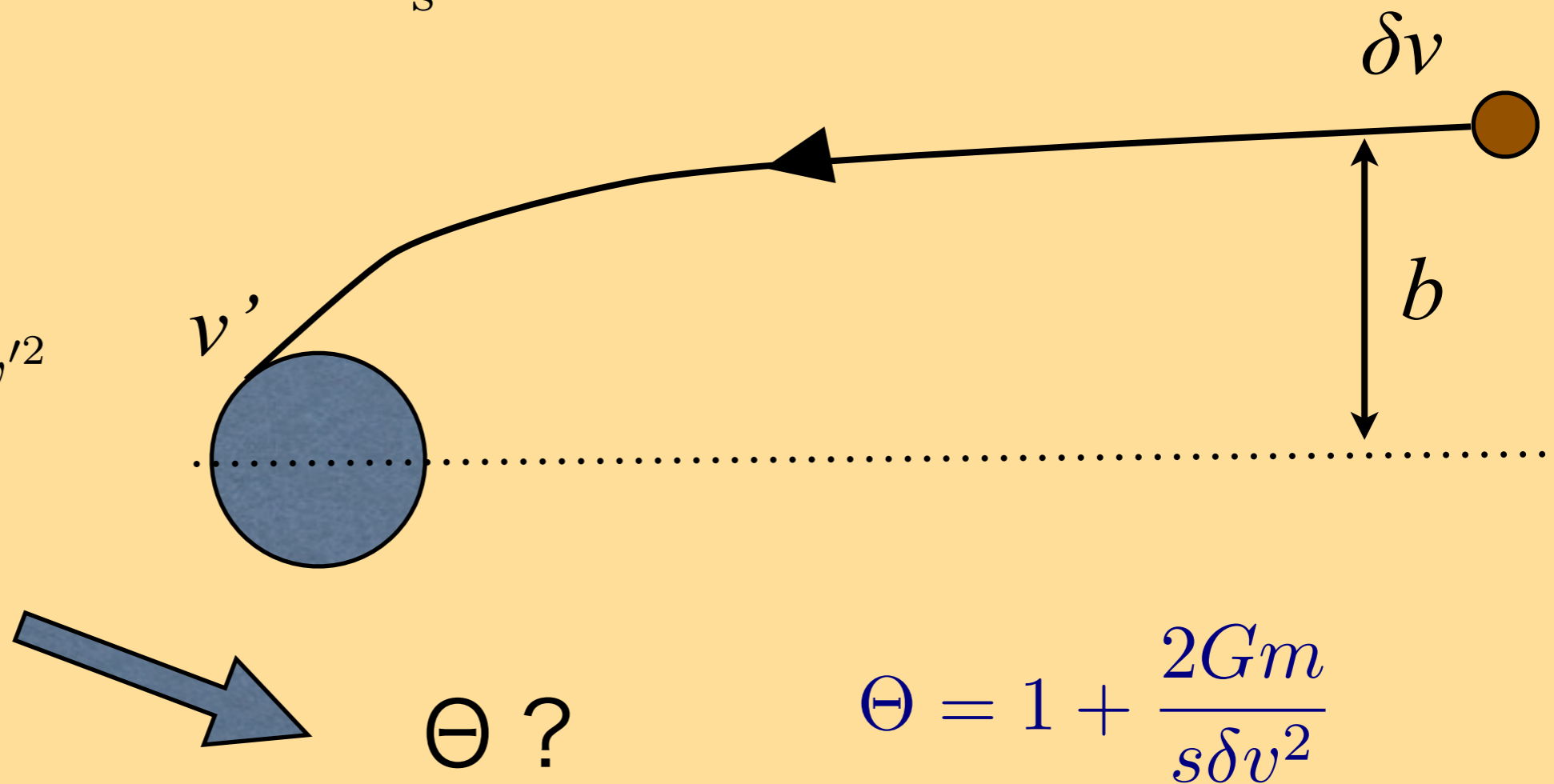
重力が効くので $\Theta > 1!$

エネルギー保存

$$\frac{1}{2} \delta v^2 = -\frac{Gm}{s} + \frac{1}{2} v'^2$$

角運動量

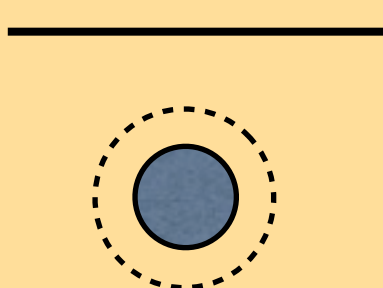
$$b \delta v = s v'$$




衝突確率

$$\frac{dm}{dt} = \sum_s \frac{\sigma \delta v}{2h_s}$$

$$\delta v = (e^2 + i^2)^{1/2} v_K, \quad h_s = ir$$

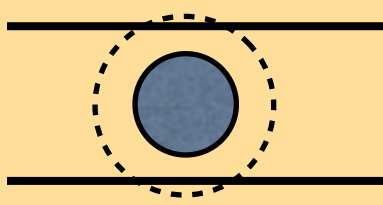


$$\frac{\sigma \delta v}{h_s} = \pi s^2 \left(1 + \frac{2Gm}{s(e^2 + i^2)v_K^2} \right) \frac{\sqrt{e^2 + i^2} v_K}{ir}$$



$$\frac{\sigma \delta v}{h_s} \simeq \pi s \left(1 + \frac{2Gm}{sr_H^2 \Omega_K^2} \right) \frac{r_H \Omega_K}{ir}$$

$$\delta v \sim r_H \Omega_K, \quad h_s = ir$$



$$\frac{\sigma \delta v}{h_s} \sim s \sqrt{1 + \frac{2Gm}{sr_H^2 \Omega_K^2}} \times h_s \times \frac{r_H \Omega_K}{h_s}$$

$$\delta v \sim r_H \Omega_K, \quad h_s < s \Theta^{1/2}$$

微惑星の成長

$$\frac{dm_1}{dt} = \sigma_{12} \frac{\Sigma_s}{2h_{s,12}} \delta v_{1,2} = r_{H,12}^2 \Sigma_s \Omega_K P_{\text{col}}$$

$$r_{H,12} = \left(\frac{m_1 + m_2}{M_*} \right)^{1/3} r$$

衝突確率 (Inaba et al. 2001)

$$\tilde{s}_{12} = \frac{s_1 + s_2}{r_{H,12}} \quad \tilde{e}^2 = \frac{(e_1^2 + e_2^2)}{r_{H,12}} \quad \tilde{i}^2 = \frac{(i_1^2 + i_2^2)}{r_{H,12}}$$

$$P_{\text{col},12} = \begin{cases} \frac{\tilde{s}_{12}^2}{2\pi} \left(17 + \frac{36}{\tilde{s}_{12} \tilde{e}_{12}^2} \right) & \tilde{e} = 2\tilde{i} \gg 1 \\ \frac{\tilde{s}_{12}^2}{4\pi \tilde{i}} \left(17 + \frac{230}{\tilde{s}_{12}} \right) & 0.2 \gtrsim \tilde{e}, \tilde{i} \gtrsim 2 \\ 11 \sqrt{\tilde{s}_{12}} & \tilde{e}, \tilde{i} \ll 1 \end{cases}$$

重力相互作用

力学的摩擦

$$m v^2 = M V^2$$

Viscus Stirring

$$\frac{dv^2}{dt} = \frac{\Sigma_s}{mh_s} \sigma_{VS} \delta v^3$$

(e.g., Ohtsuki et al. 2002)

$$\sigma_{VS} \sim \left(\frac{Gm}{\delta v^2} \right)^2 \quad h_s \sim \delta v / \Omega_K$$

$$\frac{de^2}{dt} = \Sigma_s P_{VS} \Omega_K$$

高速の場合

$$\delta v \sim \sqrt{e^2 + i^2} v_K$$

$$P_{VS,e} \sim 400 \frac{r_H^6}{r^6 e^2}$$

低速の場合

$$\delta v \sim r_H \Omega_K$$

$$P_{VS,e} \sim 73 \frac{r_H^4}{r^4}$$

暴走成長

(Wetherill & Stewart 1989; Kokubo & Ida 1996)

大きい天体の成長率

$$\frac{dM}{dt} = \frac{1}{2} \pi (S + s)^2 \left(1 + \frac{2G(M + m)}{(s + S)(v^2 + V^2)} \right) \Sigma_s \Omega_K$$

小さい天体の成長率

$$\frac{dm}{dt} = \frac{1}{2} \pi (2s)^2 \left(1 + \frac{Gm}{sv^2} \right) \Sigma_s \Omega_K$$

力学的摩擦

$$m v^2 = M V^2$$

$$\frac{T_g}{t_g} = \frac{M \dot{m}}{\dot{M} m} = \left(\frac{2s}{S + s} \right) \left(\frac{2m}{M + m} \right) \left(\frac{M}{m} \right) \left(\frac{v^2 + V^2}{2v^2} \right) \simeq 2 \left(\frac{m}{M} \right)^{1/3}$$

大きい天体ほど早く成長する。

$$M \propto \exp(t/t_{\text{col}}) \quad t_{\text{col}} = m/2\pi s^2 \Sigma_s \Omega_K$$

(e.g., Ormel et al. 2010)

寡占的成長

小さい天体が高速になると
大きい天体の成長が遅くなる。

$$v \gg \sqrt{\frac{2Gm}{s}}$$

成長率 $\frac{dM}{dt} = \pi \frac{GM S}{v^2} \Sigma_s \Omega_K$

速度進化 $\frac{dv}{dt} = \pi \left(\frac{GM}{v^2} \right)^2 v N_s \Omega_K$

大きい天体の面数密度 (Kokubo & Ida 1998)

$$N_s \approx \frac{1}{20\pi r r_H} \approx \frac{1}{20\pi r^2} \left(\frac{M}{3M_*} \right)^{-1/3}$$

$$\frac{dM}{dv} \propto \frac{v}{M^{1/3}} \rightarrow v^2 \propto M^{4/3} \rightarrow M \propto t$$

(非平衡の成長 : Kobayashi et al. 2010)

(平衡)寡占的成長

惑星による上昇

$$\frac{dv}{dt} = \pi \left(\frac{GM}{v^2} \right)^2 v N_s \Omega_K \quad N_s \approx \frac{1}{20\pi r r_H} \approx \frac{1}{20\pi r^2} \left(\frac{M}{3M_*} \right)^{-1/3}$$

ガス抵抗減衰

$$\frac{dv}{dt} = -\frac{v}{t_s} \quad t_s \propto v^{-1}$$

釣り合い $v \propto M^{1/3}$

成長率

$$\frac{dM}{dt} = \pi \frac{GMS}{v^2} \Sigma_s \Omega_K \propto M^{2/3}$$

$\rightarrow M \propto t^3$

(平衡の寡占的成長：Kokubo & Ida 2000,2002; Chambers 2006; 2008)

孤立質量

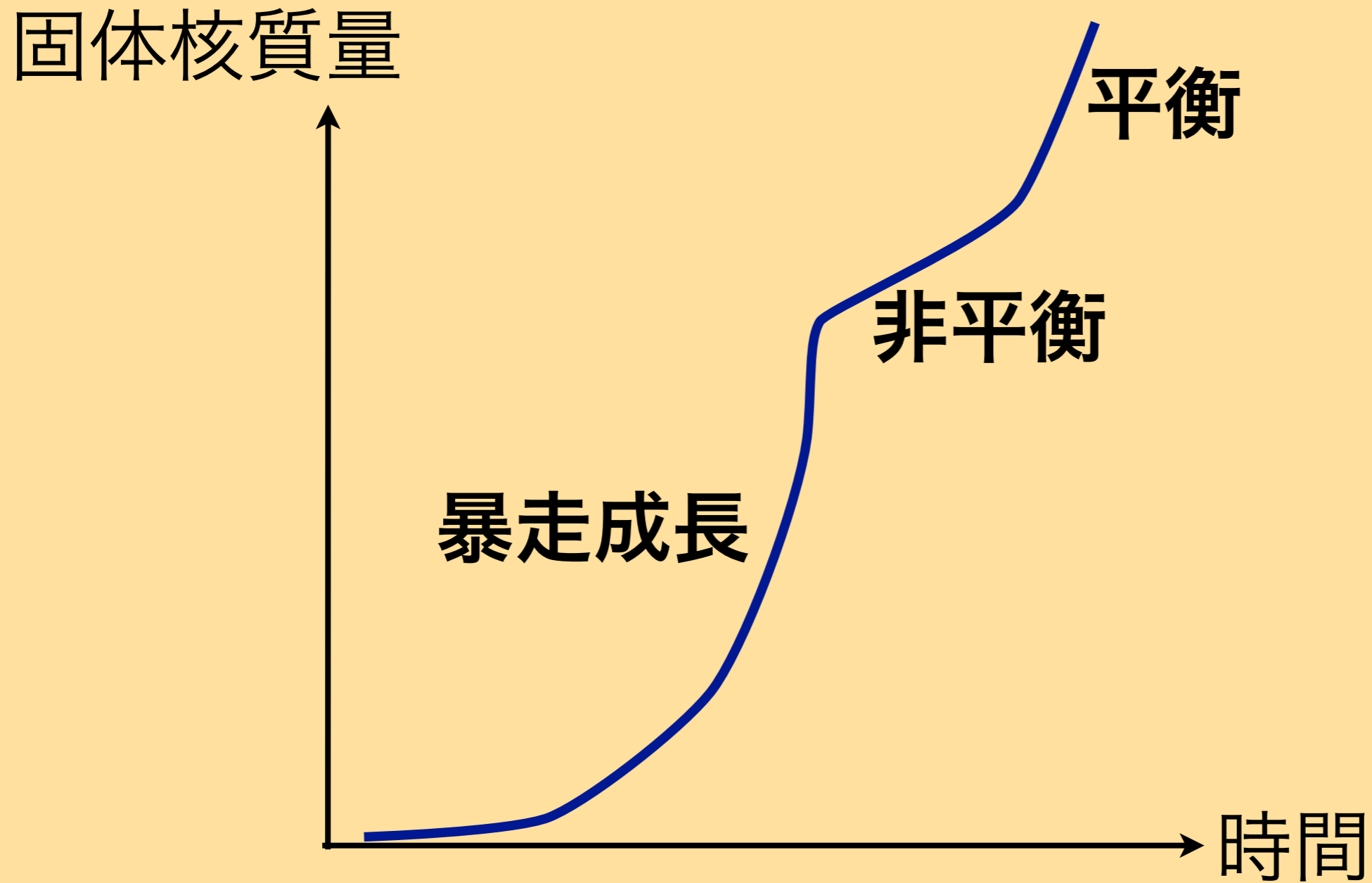
(Kokubo & Ida 2000,2002)

原始惑星は10倍の相互ヒル半径程度の間隔で並ぶ。
原始惑星がその間隔のすべての微惑星を食べたとき、
その質量は $M=20\pi r r_{HMM} \Sigma_s$ となる。

$$r_{HMM} = \left(\frac{2M}{3M_*} \right)^{1/3}$$

$$M_{\text{iso}} = 2.8 \left(\frac{\Sigma_{s,0}}{2.7 \text{ g/cm}^2} \right)^{3/2} \left(\frac{a}{5 \text{ AU}} \right)^3 \left(\frac{M_*}{M_\odot} \right)^{-1/2} M_\oplus$$

固体核の成長



果たして、このまま成長できるのだろうか？

微惑星のランダム速度

平衡の寡占的成長

$$\frac{1}{20r_{H,M}r} \left(\frac{GM}{v^2} \right)^2 v\Omega_K - 0.25\pi s^2 \rho_g v^2 / m = 0$$

原始惑星によるstirring

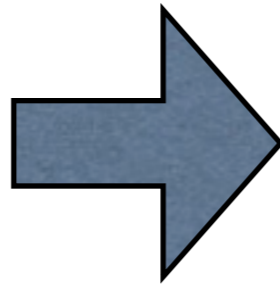
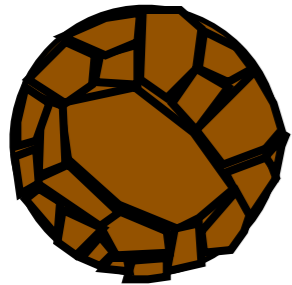
ガス抵抗

$$v = \left(\frac{9}{5\pi} \frac{m}{rs^2\rho_g} \right)^{1/5} r_{H,M}\Omega_K \propto M^{1/3}$$

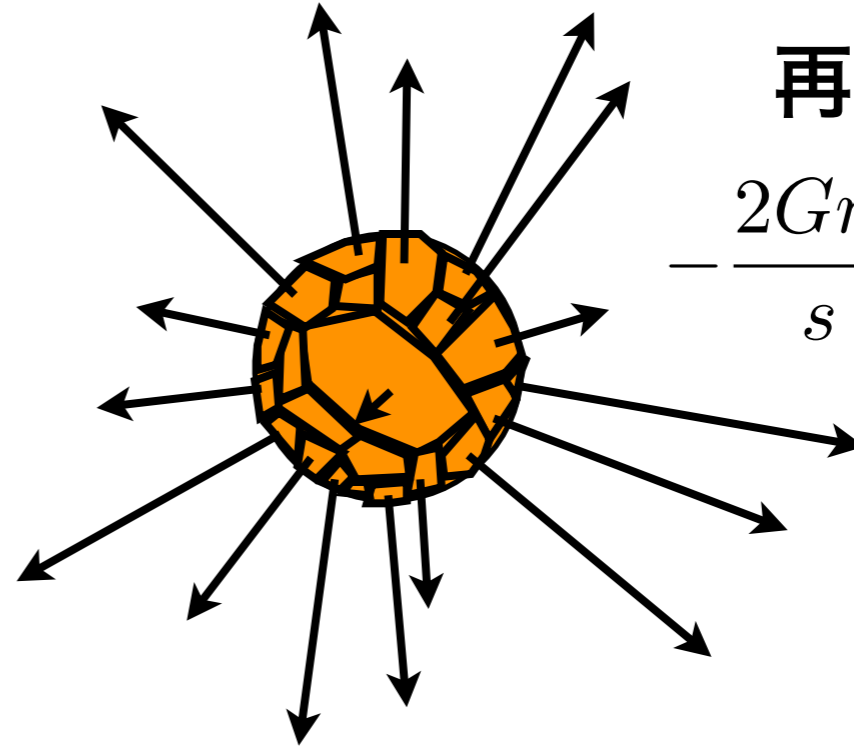
原始惑星の成長に伴い、
微惑星間のランダム速度が上がる

微惑星の破壊

$$\frac{v_r^2}{2} = -\frac{2Gm}{s} + \frac{v_{col}^2}{2}$$



$$v_{eject} = \epsilon v_{col}$$



再集積のために

$$-\frac{2Gm}{s} + \frac{v_{eject}^2}{2} < 0$$

再集積条件

$$v_r < \sqrt{\frac{1 - \epsilon^2}{\epsilon^2}} \sqrt{\frac{2Gm}{s}}$$

脱出速度の3倍程の
相対速度で破壊

微惑星の破壊

速度の比較

$$\frac{r_{H,M} \Omega_K}{v_{esc}} \approx 10 \left(\frac{M}{0.1 M_{\oplus}} \right)^{1/3} \left(\frac{r}{1 \text{ AU}} \right)^{1/2} \left(\frac{s}{10 \text{ km}} \right)^{-1}$$

火星質量程度でも破壊が起こる！

破壊を考慮した固体核形成モデルの構築の必要性！

- 微惑星の破壊による減少
- 破片の原始惑星への集積

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