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惑星移動 (Type I migration)

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固有振動数 = 会合周波数の整数倍 $r_L = \mathop{\mathsf{C}}\limits_{\mathsf{e}}^{\mathfrak{A}} 1 \pm \frac{1}{m\mathfrak{G}}^{\mathfrak{G}^{2/3}} r_p$

on model would require an extremely small intrinsic viscosity to account for the gap size. calculations on gap formation, it was often assumed that the waves excited near LRs are loca disk, the effective viscosity may not be adequate to inhibit wave propagation. If the per small and amplitude of the wave is linear, the excited density waves can propagate to large d d transfer angular momentum effectively over a large region of the disk (see Fig. 1b). The pro sture of the disk at the location where the waves are dissipated. Thus, the width of the gap may t derived on the basis of local dissipation. This idea has been suggested as a mechanism for inc division in Saturn's rings by Goldreich & Tremaine (1978). Nonlocal dissipation of density ed in Saturn's rings (Cuzzi, Lissauer, & Shu 1981).

, we investigate up he feet al why oppop agation on the process of gap formation. We assume the day an embedded protoplanet is small. In this limit, wave propagation and viscous dissipation





エピサイクル運動 (周転円運動) 周期:*T_K*

 $v = aeW_{K} = V_{K}e$

Tanaka et al. 2002

- ・等温ガス円盤, 3D構造
- ・トルク

$$t_{I} = \frac{1}{2.7 + 1.1q} \left(\frac{M}{M_{*}}\right)^{-1} \left(\frac{S_{g}r^{2}}{M_{*}}\right)^{-1} \left(\frac{C_{s}}{V_{K}}\right)^{2} W_{K}^{-1} \qquad S_{g} \mu r^{-q}$$

$$> 5 \quad 10^4 \mathring{\mathcal{C}} \frac{M}{M_{\text{Earth}}} \overset{\ddot{0}^{-1}}{\overset{\circ}{\sigma}} \overset{\mathcal{R}}{\underset{e}{\circ}} \frac{r}{1 \text{ AU}} \overset{\ddot{0}^{3/2}}{\overset{\circ}{\sigma}} \text{ yr}$$

「惑星落下問題」

Paardekooper & Mellema 2006

- ・非等温ガス円盤
- 3D流体計算
 with 輻射エネルギー輸送
- ・トルク



Fig. 1. Total torque on a 5 M_{\oplus} planet as a function of time for three different midplane densities, together with the isothermal result. The torques are normalized to the analytical value found by Tanaka et al. (2002), which is reproduced by the isothermal simulation. For high densities (and thereby for high opacities) the torque becomes positive, indicating outward migration.



Type I migration in optically thick accretion discs

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孤立質量





原始惑星系円盤の質量



Beckwith et al. 1990, AJ 99, 924



Ida & Lin 2004



FIG. 9.—Theoretically predicted distribution based on the core accretion model for gas giant planets (for the range of parameters we used, see text). Cores are truncated by $M_{c,iso}$. Gas accretion is truncated by $(a) M_{g,iso}$, $(b \text{ and } c) M_{g,ih}$, and $(d) M_{g,vis}$. We adopt $\Delta a_g = 2r_H$ in (a), the critical Hill radius $r_{H,c}$ being h and 1.5h in (b) and (c), respectively, and $\alpha = 10^{-3}$ in (d). The green filled circles and the blue crosses represent rocky and icy planets with gaseous envelopes less massive than their cores. The green and blue open circles represent gas-rich rocky and icy planets with gaseous envelopes that are 1–10 times more massive than their cores. The red filled circles are present gas giants with envelopes more massive than 10 times their cores. For comparison, we also plot observational data of extrasolar planets in (d). The dashed ascending lines correspond to radial velocity amplitude of 100 (*upper line*), 10 (*middle line*), and 1 m s⁻¹ (*lower line*), assuming that the host star mass is 1 M_{\odot} .



太陽系の形成に特化したモデル:

The Grand Tack Scenario (Morbidelli et al 2012)



Walsh et al. 2011



・小惑星





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